

Extragalactic Jets: Theory and Observation from Radio to Gamma Ray
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The Case of the 300 kpc Long X-ray Jet in PKS 1127-145 at $z=1.18$

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Abstract.

The complex X-ray morphology of the 300 kpc long X-ray jet in PKS 1127-145 ($z=1.18$ quasar) is clearly discerned in a ~ 100 ksec *Chandra* observation. The jet X-ray surface brightness gradually decreases by an order of magnitude going out from the core. The X-ray spectrum of the inner jet is relatively flat with $\alpha_X = 0.66^{+0.15}_{-0.15}$ and steep in the outer jet with $\alpha_X = 1.0^{+0.2}_{-0.2}$. The X-ray and radio jet intensity profiles are strikingly different, with the radio emission peaking strongly at the two outer knots while the X-ray emission is strongest in the inner jet region. We discuss the constraints implied by these data on the X-ray emission models and conclude that “one-zone” models fail and that at least a two-component model is needed to explain the jet’s broadband emission. We propose that the X-ray emission originates in the jet proper while the bulk of the radio emission comes from a surrounding jet sheath. We also consider intermittent jet activity as a possible cause of the observed jet morphology.

1. Introduction

The origin of the X-ray emission in quasar jets is puzzling (see Harris & Krawczynski 2006, for review) because a straight extrapolation of the synchrotron radio-to-optical continuum into the X-rays underpredicts the luminosity of powerful *Chandra* large scale jets. Thus, the same single power-law population of electrons cannot produce the radio, optical, and X-ray emission in the framework of a homogeneous one-emission zone approximation (e.g. Sambruna et al. 2004). The synchrotron self-Compton (SSC) process cannot easily explain the data

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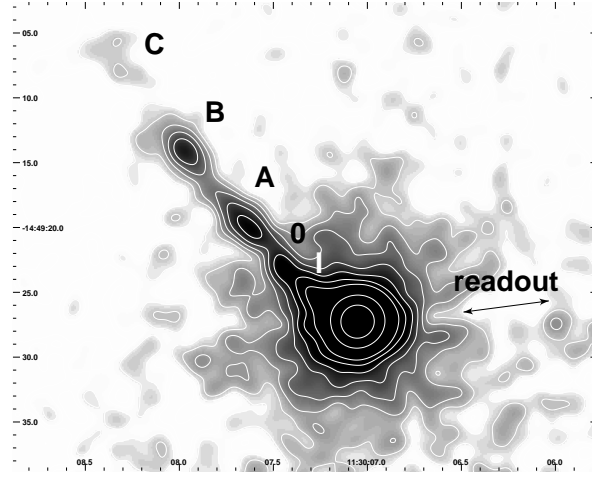


Figure 1. *Chandra* ACIS-S exposure corrected image of PKS1127-145 ($E=0.3-7$ keV) and smoothed with the Gaussian kernel ($1\sigma=0.615$ arc-sec). Contours are factor of two in log scale starting at 10^{-10} photons $\text{cm}^{-2} \text{s}^{-1} \text{pixel}^2$. The two final contours within the core are at 10^{-6} and 10^{-7} photons $\text{cm}^{-2} \text{s}^{-1} \text{pixel}^2$ (1 pixel= $0.123''$).

because it does not produce enough X-rays at the equipartition fields and therefore requires large departures from the minimum-power condition (Chartas et al. 2000; Harris & Krawczynski 2002; Kataoka & Stawarz 2005). The alternative scenario involves inverse-Compton scattering of the Cosmic Microwave Background (IC/CMB) photons implying large jet bulk Lorentz factors (Γ) at hundreds of kpc from the active nuclei (Tavecchio et al. 2000; Celotti et al. 2001; Schwartz 2002). In this paper we discuss the constraints on the X-ray emission models given by the *Chandra* and VLA observations of the PKS 1127-145 jet.

2. Jet Properties

One of the longest X-ray jets known (see <http://hea-www.harvard.edu/XJET/>) is associated with the redshift $z=1.18$ radio-loud quasar PKS 1127-145 (Fig. 1). This jet poses several questions for X-ray emission models (see Siemiginowska et al. 2002, 2007). The jet morphology is complex with several knots seen in the X-ray image (Fig. 1) with the furthest knot C being the weakest. The prominent knots A and B are connected by continuous X-ray emission that stops beyond knot B at ~ 22 arcsec from the core. The VLA radio maps show low brightness emission along the jet (Fig. 2). The one-sided jet shows an X-ray surface brightness which declines with the distance from the core, while the radio brightness increases (Fig. 3[left]). The X-ray to radio intensity ratio decreases along the jet. Fig.3[right] shows the changes to spectral index along the jet. The radio spectrum is steep and X-ray spectrum is flat in the inner jet regions. There are only upper limits to the knots optical brightness, but they are too high for constraining the emission models. Thus only radio and X-ray observations can

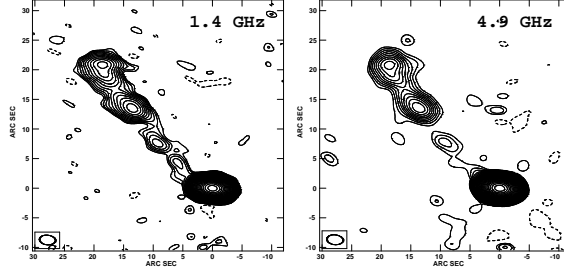


Figure 2. VLA 1.4 and 4.9 GHz images of PKS 1127–145. The lowest contour levels plotted are $0.75 \text{ mJy beam}^{-1}$ (1.4 GHz) and $0.39 \text{ mJy beam}^{-1}$ (4.9 GHz) and increase by factors of $\sqrt{2}$ up to peaks of 5.46 and 4.14 Jy/beam, respectively. The (uniform weighted) beam sizes are: $2.71'' \times 1.62''$ at $\text{PA}=84.4^\circ$, and $2.77'' \times 1.66''$ at $\text{PA}=84.0^\circ$, respectively.

be used to study emission processes in this jet. We discuss the constraints on the models below.

3. Jet Emission Models

A detailed discussion of the model constraints can be found in Siemiginowska et al. (2007). Here we give a summary of the main results:

3.1. One Zone Model Failures

Synchrotron X-ray emission requires a continuous acceleration of high energy particles along the jet. The expected distribution of the freshly accelerated electrons gives $\alpha_X = 0.5$ as observed in the inner jet, but the radio spectrum is steep with $\alpha_R > 1.0$ (Fig.3). A different cooling mechanism or difference in electron populations emitting radio and X-rays is needed.

IC/CMB models in the simplest version involve an adiabatically expanding jet. In this case the observed inverse-Compton (X-rays) and synchrotron (radio) intensity ratio is different than the predicted one. In the case of an efficient jet confinement with some deceleration and B-field amplification it is possible to get an agreement between the predicted and the observed luminosity ratios, but the required B-field amplification is high $[B(r_C)/B(r_o)]^2 \sim 50$. In this case for the equipartition at r_o there will be no equipartition at r_C . The expected radio spectrum should steepen more significantly than the X-ray spectrum. The observed behavior is different: α_R decreases, while α_X increases (Fig.3[right]).

3.2. Two-Component Model: The Jet and a Sheath

Radio and X-ray emissions are produced in two separate regions. The X-ray emission comes from the proper relativistically moving jet and the radio emission from the sheath - a slow moving radial extension of the jet boundary layer. The X-ray emission is due to synchrotron or IC/CMB process.

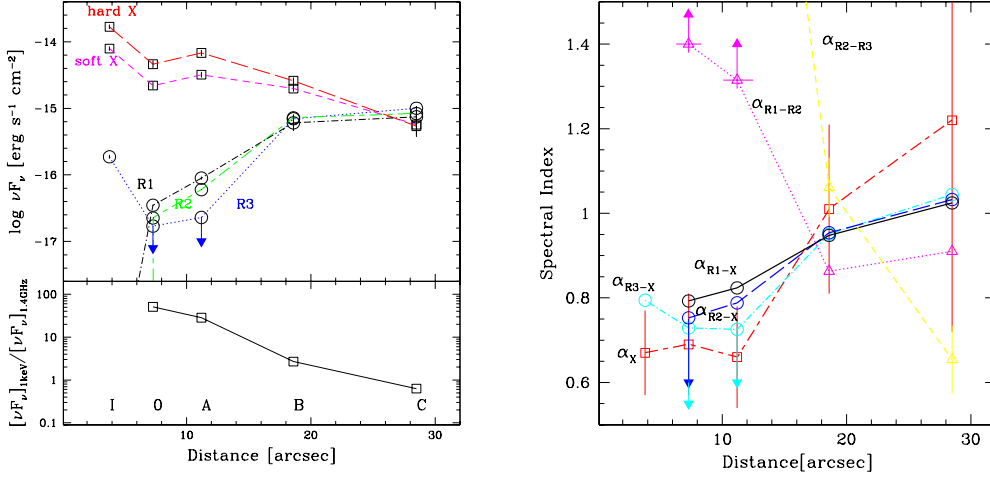


Figure 3. **Left:** Flux along the jet for each observed frequency: long dashed line – hard X-rays (2-10 keV); short dashed line – soft X-rays (0.5-2 keV); dotted line – R1 = 1.4 GHz; short-long-dashed – R2 = 5 GHz; and dash-dot – R3 = 8.5 GHz. Lower panel shows the flux ratio (1keV/1.4GHz) along the jet. **Right:** Spectral index ($S_\nu \sim \nu^{-\alpha}$) along the jet: α_X – X-ray squares, and long-short dash; α_{R3-X} – 8.5 GHz to 1 keV circles, dot-dashed line; α_{R2-X} – 5 GHz to 1 keV, circles, dashed line; α_{R1-X} – 1.4 GHz to 1keV, circles, solid line; α_{R1-R2} – 1.4 GHz to 5 GHz, triangles, dotted line, α_{R2-R3} – 5 GHz to 8.5 GHz, triangles dashed line.

Proper Jet: The observed smooth decrease of the X-ray flux, when unconnected with the increase of the radio one, is expected as a result of the decrease in the inverse-Compton or synchrotron luminosities along decelerating and/or expanding jet.

Sheath: The sheath dominates the jet’s radio emission and has two cooling regions: (1) radiative cooling at the inner sheath: $t_{rad} < 10 \text{ Myr} < t_{dyn} = 30R_{10} \text{ Myr}$. Due to the frequency dependent losses a steep radio spectrum is expected in this regime. (2) Adiabatic cooling in the outer sheath: $t_{rad} > t_{dyn} < 10R_{10}\text{Myr}$. Because of the frequency independent losses a flat spectrum is expected.

3.3. Modulated Jet Activity?

Can the morphology of the >300 kpc long jet be a result of the intermittent activity of the quasar? The knots are potentially too large (>10 kpc) to be considered a result of extended shock waves formed within a continuous jet outflow. They might form during episodes of separate continuous activity. Then the size gives a duration for each epoch of $\sim 10^5$ years. The radio core can be related to blazar phenomena (Błażejowski et al. 2004), but the core also resembles a compact GPS object at this time.

4. Summary

The key results of this project:

- The jet is long: radio and X-ray jet emission is detected up to ~ 300 kpc (projected distance) from the core.
- The jet X-ray brightness is decreasing while the radio brightness is increasing with distance from the quasar.
- The one component X-ray/radio emission models failed to explain the observations.
- X-ray and radio jet properties suggests two separate components for the jet emission.
- Intermittent quasar activity might be reflected in the observed jet morphology: each knot represents a continuous jet activity with a duration of 10^5 years.

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References

- Błażejowski, M. et al. 2004, ApJ, 600, L27
 Celotti, A., Ghisellini, G., & Chiaberge, M. 2001, MNRAS, 321, L1
 Chartas, G., et al. 2000, ApJ, 542, 655
 Cheung, C. C., Stawarz, L., & Siemiginowska, A. 2006, ApJ, 650, 679
 Harris, D. E., & Krawczynski, H. 2002, ApJ, 565, 244
 Harris, D. E., & Krawczynski, H. 2006, ARA&A, 44, 463
 Kataoka, J., & Stawarz, L. 2005, ApJ, 622, 797
 Sambruna, R. M. et al. 2004, ApJ, 608, 698
 Schwartz, D. A. 2002, ApJ, 569, L23
 Siemiginowska, A. et al 2002, ApJ, 570, 543
 Siemiginowska, A. et al. 2007, ApJ, 657, 145
 Tavecchio, F., Maraschi, L., Sambruna, R. M., & Urry, C. M. 2000, ApJ, 544, L23